# TR/04/86

March 1986

Global extrapolation procedures for special and general initial value problems.

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ABSTRACT

Two- and three- grid global extrapolation procedures are considered for the special and general initial value problems of arbitrary order q. Extrapolation formulas are developed for consistent numerical methods of arbitrary order p .

The global extrapolations of a number of existing numerical methods are considered and tested on three problems from the literature.

#### 1. THE SPECIAL INITIAL VALUE PROBLEM

#### 1.1 Introduction

Consider the special initial value problem of order q given by (1)  $y^{(q)}(t) = f(t, y)$ ;  $y^{(r)}(t_0) = z_r$ ,  $r = 0, 1, ..., q^{-1}$ and suppose that the solution is sought at time  $t = T < \infty$ .

The interval of integration will be divided, first of all, into N sub-intervals each of width h so that  $Nh = T - t_0$ , giving a discretization or grid G<sub>1</sub> consisting of the N+1 points  $t_{n,1} = t_0 + nh$ (n = 0, 1, ..., N). The theoretical solution of (1) at  $t = t_{n,1}$  is clearly  $y(t_{n,1})$  and the notation  $y_{n,1}$  will be used to denote the solution of an approximating method at the same point  $t_{n+1}$  of  $t_{n,1}$  of G<sub>1</sub> (n = 0, 1, ..., N).

The application of a convergent numerical method M to find the solution yields, at the point  $T = t_{N,1}$  of G<sub>1</sub> the magnified error function (Lapidus and Seinfeld [5; p.242], Henrici [3; p.80]) or the global error (Verwer and de Vries [10]) in the form

(2)  $\varepsilon_{N,1} = C_{p+q} h^p y^{(p+q)} (T) + c_{p+q+2} h^{p+2} y^{(p+q+2)} (T) + c_{p+q+4} h^{p+4} y^{(p+4)} y^{(p+q+4)} (T) + ....,$ where  $p \ge 1$  is the *order* of the numerical method and  $C_{p+q}$  is its *error constant*. The term in  $h^p$  in (2) gives the time component of the principal part of the local truncation error when M is associated with the solution of a time-dependent partial differential equation.

#### 1.2 Global extrapolation using two discretizations

Suppose now that the interval of integration is divided into 2N subintervals each of width  $\frac{1}{2}h$  giving a discretization G<sub>2</sub> consisting of the 2N+1 points  $t_{i,2} = t_0 + \frac{1}{2}$  ih (i = 0, 1,..., 2N). Clearly the points  $t_{r,2}$ (r=0,2,4,...,2N) of G<sub>2</sub> are coincident with the points  $t_{n,1} = t_0 + nh$ (n = 1,2,...,N) of G<sub>1</sub>. The notation  $y_{i,2}$  will be used to denote the

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solution of the method M at the points  $t_{i,2} \ (i=0,1 \ , \ldots , 2N)$  of the grid  $G_2$ 

The application of M to find the solution at the point  $T = t_{2N,2}$ of G<sub>2</sub> generates the global error

(3) 
$$\epsilon_{2N,2} = 2^{-p} c_{p+q} h^{p} y^{(p+q)} (T) + 2^{-p-2} c_{p+q+2} h^{p+2} y^{(p+q+2)} (T)$$
$$+ 2^{-p-4} c_{p+q+4} h^{p+4} y^{(p+q+4)} (T) + \dots$$

which, like  $\varepsilon_{N,1}$  is  $O(h)^p$  so that  $y_{N,1}$  and  $y_{2N,2}$  are both approximations of order p to y(T).

(4) Consider, now, the approximation  $y^{(E)} = \alpha y_{2N2} + (1-\alpha)y_{N1}$ 

and the associated global error

(5) 
$$\varepsilon^{(E)} = \alpha \varepsilon_{2N,2} + (1 - \alpha) \varepsilon_{N,1}.$$

It is easy to show that the term in  $h^p$  in (5) vanishes when the parameter a takes the value

(6) 
$$\alpha = 2^{p}/(2^{p}-1)$$
 with  $1-\alpha = 1/(1-2^{p})$ .

The global extrapolation carried out using the two discretizations  $G_1$  and  $G_2$  has thus produced an approximation  $y^{(E)}$  defined by (4), which is of order p+2 provided  $\alpha$  takes the value given by (6).

## 1.3 Numerical results

The global extrapolation procedure described in §1.2 was tested on the following problem

Problem 1. This problem is given by

$$y'_1 = y_2$$
;  $y_1(0) = 1$ ,  
 $y'_2 = -\frac{y_2}{t} + y_1^3 - 3y_1^5$ ;  $y_2(0) = 0$ .

The problem has a singularity in  $y'_2$ ; therefore, a fully implicit method must be used to obtain the solution. Noting that the problem has the vector form  $\underline{y'}(t) = \underline{f}(t, \underline{y})$ , the solution was obtained using the first order backward Euler method

$$y(t+h) - \ell f(t+h, y(t+h)) = y(t) ,$$

the Newton-Raphson method for an algebraic system (of order 2) being employed to compute  $\underline{y}(t+h)$ .

The problem has theoretical solution

$$y_2(t) = (1+t^2)^{-\frac{1}{2}}$$
,  $y_2(t) = -t(1+t^2)^{-\frac{3}{2}}$ 

and the maximum error moduli at time t = 0.25 are given in Table 1. It may be seen from Table 1 that the errors relating to one grid are decreasing by a factor of 2 (approximately) as h is successively halved, while the errors following global extrapolation of the solution are decreasing by a factor of 4 (approximately).

#### 1.4 Global extrapolation using three discretizations

Suppose, finally, that the interval of integration is divided into 3N subintervals each of width  $\frac{1}{3}$ h giving a discretization G<sub>3</sub> consisting of the 3N+1 points  $t_{j,3} = t_0 + \frac{1}{3}$  jh (j=0,1,...,3N). It is clear that the points  $t_{s,3}$  (s = 0,3,6,...3N) of G<sub>3</sub> are coincident with the points  $t_{n,1}$  (n = 0,1,...,N) of the original grid G<sub>1</sub>. The notation  $y_{j,3}$  will be used to denote the solution of the numerical method M at the points  $t_{j,3}$  (j = 0,1,...,3N) of G<sub>3</sub>.

The application of M to find the solution at the point  $T = t_{3N,3}$ of G<sub>3</sub> gives a third approximation to y(T) and generates the global error

(7) 
$$\epsilon_{3N,3} = 3^{-p} c_{p+q} h^{p} y^{(p+q)}(T) + 3^{-p-2} c_{p+q+2} y^{(p+q+2)}(T) + 3^{-p-4} c_{p+q+4} y^{(p+q+4)}(T) + \dots ,$$

Which, like  $\varepsilon_{N,1}$  and  $\varepsilon_{2N,2}$  is  $0(h^p)$ , so that the third approximation to y(T), given by  $y_{3N,3}$ , is also of order p.

Considering the approximation

(8) 
$$y^{(E)} = \alpha y_{3N,3} + \beta y_{2N,2} + (1 - \alpha - \beta) y_{N,1}$$

and the resulting global error

(9) 
$$\varepsilon^{(E)} = \alpha \varepsilon_{3N,3} + \beta \varepsilon_{2N,2} + (1 - \alpha - \beta) \varepsilon_{N,1}$$

where  $\alpha$  and  $\beta$  are parameters, it may be shown that the terms in  $h^p$  and  $h^{p+2}$  in (9) vanish when

(10) 
$$\alpha = 3^{p+3} / (5 + 3^{p+3} - 2^{p+5}), \beta = -2^{p+5} / (5 + 3^{p+3} - 2^{p+5})$$

with, consequently,  $1-\alpha-\beta = 5/(5+3^{p+3}-2^{p+5})$ .

This global extrapolation, which uses the three discretizations  $G_1$ ,  $G_2$  and  $G_3$ , has produced an approximation  $y^{(E)}$  defined by (8) which is of order p+4 provided  $\alpha$  and  $\beta$  take the values in (10).

## 1.5 Numerical results using the three-grid extrapolation.

*Problem 2* (Stiefel and Bettis [7]). This is the "almost periodic" problem given by

 $z''(t) + z(t) = 0.001e^{it}$ ; z(0) = 1, z'(0) = 0.9995i,  $z(t) \in \mathbb{C}$ . The analytic solution of this problem is given by

$$\begin{split} u(t) &= \cos t \, + \, 0.0005t \sin t \, , u \in \, \mathrm{I\!R} \, , \\ v(t) &= \sin t \, - \, 0.0005t \cos t \, , v \in \, \mathrm{I\!R} \, , \\ z(t) &= u(t) \, + \, iv(t) \end{split}$$

and represents the motion of the point z(t) on a perturbation of a circular orbit. The distance of this point from the centre of the orbit at time t is given by  $\gamma(t) = \{u^2(t)+v2(t)\}^{\frac{1}{2}}$  and the error modulus  $|E(\gamma)|$  of the computed value of  $\gamma$  was determined at  $t = 40 \pi$  for the three-discretization extrapolations of the fourth order method based on the (2,2) padé approximant of Twizell and Khaliq [8], the fourth order method of Cash [1] and the sixth order method of Cash [1].

The computed results were found using the time steps  $h = \frac{\pi}{4}, \frac{\pi}{5}, \frac{\pi}{6}, \frac{\pi}{9}$  and  $\frac{\pi}{12}$  (on the first discretization G<sub>1</sub>); the results for the two fourth order methods are given in Table 2 and for Cash's sixth order method in Table 3. The solution at the first time step for every numerical experiment was computed using the second order Taylor series approximation.

It is not the purpose of the present paper to compare the relative merits of the methods developed in [1,8]. The aim of the paper is to show that global extrapolation, as detailed above, increases the order of the method being used. It is clear from the two tables that this aim has been achieved for all three methods tested on Problem 2. *Problem 3* (Van Dooren [9]). This is the nonlinear Duffing equation

 $y''(t) + y(t) + y^{3}(t) = F \cos \Omega t$ ; y(0) = A, y'(0) = 0with F = 0.002,  $\Omega$ =1.01 and A = 0.200426728067 (Chawla and Rao [2]).

This problem was solved using the second order P-stable method based on the (1,1) padé approximant [8] and the more accurate second order method based on the (1,2) padé approximant [8] which, using the notation of Lambert and Watson [4], has periodicity interval  $H^2 \in (0,7.2)$ .

The global extrapolation procedure using three discretizations was also carried out, increasing the orders of each of the two numerical methods to six. The step size h was given the values  $\pi/5, \pi/10, \pi/20$ , and  $\pi/40$  and the solution computed at time t = 40 $\pi$ . Van Dooren [9] gives the solution of Problem 3 in the form

(16) 
$$y(t) = \sum_{i=0}^{4} a_{2i+1} \cos[(2i+1)\Omega t],$$

where

$$a_1 = 0.200179477536 , \qquad a_3 = 0.000246946143 , \\ a_5 = 0.000000304014 , \qquad a_4 = 0.00000000374 , a_5 = 0.0 ,$$

noting that the order of (16) is nine, with a precision of the coefficients of  $10^{-12}$ .

The errors using the two numerical methods with one and three grids are given in Tabel 4 where, for comparison purposes, the results of chawla and Rao [2] relating to their sixth order method  $M_6(0)$  are reproduced,

It is seen from Table 4 that the order of the two methods based on the (1,1) and (1,2) padé approximants [8] are duly increased by the global extrapolation procedure of §1.4. The results using the (1,1) method with three discretizations are inferior to those of Chawla and Rao [2], while those using the (1,2) method with three discretizations are better than those in [2]. particularly for the value  $h = \pi/5$ .

#### 2. THE GENERAL INITIAL VALUE PROBLEM

Consider now the general initial value problem of order q given by (11)  $y^{(q)}(t) = f(t, y(t), y'(t), y''(t), ..., y^{(q-1)}(t))$ ;  $y^{(r)}(t_0) = z_r, r = 0, 1, ..., q-1$ 

and suppose, again, that the solution is sought at time  $t = T < \infty$ .

As in §1, the interval of integration will be divided into subintervals in three ways giving rise to the same three grids  $G_1, G_2, G_3$ .

The application of a convergent numerical method M to find two approximations  $y_{N,1}$  and  $y_{2N,2}$  yields the global errors

(12) 
$$\varepsilon_{N,1} = c_{p+q} h^{p+q} y^{(p+q)}(T) + c_{p+q+1} h^{(p+q+1)} y^{(p+q+1)}(T) + c_{p+q+2} h^{p+q+2} y^{(p+q+2)}(T) +$$
  
and

(13) 
$$\epsilon_{2N,2} = 2^{-p} c_{p+q} h^{p} y^{(p+q)}(T) + 2^{-p-1} c_{p+q+1} h^{p+q+1} y^{(p+q+1)}(T) + 2^{-p-2} c_{p+q+2} h^{p+q+2} y^{(p+q+2)}(T) + \dots$$

respectively, where, as in §1, p is the order of M.

Considering, now, the approximation  $y^{(E)}$  of (4) and the associated global error given by (5), it is easy to show that the term in  $h^p$  in (5) vanishes when the parameter  $\alpha$  takes the value given by (6). The global extrapolation procedure involving the two discretizations G<sub>1</sub> and G<sub>2</sub> described in §1.2 for the special initial value problem (1), is therefore valid for the general initial value problem (11) but, this time, the order of the extrapolation is only p+1.

The application of M on the third grid  $G_3$  generates the approximation  $y_{3N_3}$  and the associated global error function

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(14) 
$$\epsilon_{3N,3} = 3^{-p} c_{p+q} h^{p} y^{(p+q)}(T) + 3^{-p-1} c_{p+q+1} h^{p+q+\overline{1}} y^{(p+q+1)}(T) + 3^{-p-2} c_{p+q+2} h^{p+q+2} y^{(p+q+2)}(T) + 3^{-p-2} c_{p+2} h^{p+q+2} h^{p+q$$

Considering the approximation  $y^{(E)}$  of (8) and the associated global error given by (9), it may be shown that the terms in and  $h^p h^{p+1}$  and in (9) vanish when the parameters  $\alpha$  and  $\beta$  take the values

(15) 
$$\alpha = 3^{p+1}/(1+3^{p+1}-2^{p+2})$$
 and  $\beta = -2^{p+2}/(1+3^{p+1}-2^{p+2})$ 

so that  $1-\alpha-\beta = 1/(1+3^{p+1}-2^{p+2})$ . The three-discretization extrapolation

of the general initial value problem (11) is thus only of order p+2, compared with p+4 for the special problem (1).

Putting p = q = 1 and using the two grids  $G_1$  and  $G_2$  gives the second order global extrapolation procedure for parabolic equations outlined by Verwer and de Vries [10], while using all three grids  $G_1$ ,  $G_2$  and  $G_3$  with p = q = 1 gives the third order algorithm of those authors. Putting p = 2and q = 1 gives the values of the parameters  $\alpha$  and  $\beta$  associated with the stable, two- or three-discretization global extrapolations of the wellknown Crank-Nicolson method.

It may also be seen that, putting p=q=N=1 in the above analysis, and using  $G_1$  and  $G_2$ , gives the parameter  $\alpha$  for the L<sub>0</sub>-stable second order *local extrapolation* method of Lawson and Morris [6] which was based on the well-known fully implicit first order method for parabolic equations. The local extrapolation of the well-known Crank-Nicolson method for parabolic equations is also described by the above procedure (using  $G_1$  and  $G_2$ ) with q=N=1 and p=2; this local extrapolation has a stability restriction.

Table 1.Maximum error moduli for Problem 1 using the first order<br/>fully implicit method.

Grids	1	2
Order	1	2
h=1/16	0.56E-2	0.25E-3
h=1/32	0.29E-2	0.62E-4
h=1/64	0.15E-2	0.15E-4
h=1/128	0.76E-3	0.38E-5

Table 2.Error moduli for problem 2 using two fourth order methods [1, 8]

Method	(2,2)padé [8]		Cash[1]		
Girds	1 3		1	3	
Order	4	4 8		8	
h=π/4	0.43E-2	0.21E-5	0.46E-2	0.25E-5	
h=π/5	0.13E-2	0.32E-6	0.14E-2	0.38E-6	
h=π/6	0.53E-3	0.63E-7	0.56E-3	0.76E-7	
h=π/9	0.88E-4	-0.85E-9	0.94E-4	0.13E-8	
h=π/12	0.27E-4	0.79E-10	0.28E-4	0.52E-10	

Grids	1	3
Order	6	10
h=π/4	0.24E-4	0.11E-8
$h=\pi/5$	0.49E-5	0.14E-9
h=π/6	0.14E-5	0.26E-10
h=π/9	0.11E-6	0.62E-12
h=π/12	0.18E-7	0.46E-13

Table 3.	Error moduli for Problem 2 using the sixth order
	method of Cash [1]

Table 4. Error moduli for problem 3

Method	(1,1)padé [8]		(1,1)padé [8]		Chawla and Rao [2]
Girds	1	3	1	3	1
Order	2	6	2	6	6
h=π/5	0.36E-1	0.20E-1	0.87E-1	0.77E-4	0.14E-2
h=π/10	0.14	0.28E-1	0.23E-1	0.12E-5	0.22E-4
h=π/20	0.35E-1	0.41E-5	0.58E-2	0.19E-7	0.34E-6
h=π/40	0.87E-1	0.54E-7	0.15E-2	0.52E-8	0.54E-8

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